



# Physics with high energy photons

*(The **quicksilver** bullet 😊)*

PHENIX Focus – May 30, 2006

G. David, BNL

# Overview – and rules of the game



The beauty and curse of photons, as **penetrating probes**

The basic (pQCD) processes – **colinearity**

Additional processes **in the medium**

Experiments (a not too random selection – includes **controversial** stuff, and **not** meant to be a survey of most recent results)

Recapitulation: **discriminating** between different processes with photons

This year's PHENIX-Focus tradition: talk about context, history

I won't always let reality get in my way ☺

- just because I talk about an interesting signal it doesn't necessarily mean the measurement is feasible at RHIC/PHENIX
- and if it is, often it is awfully hard

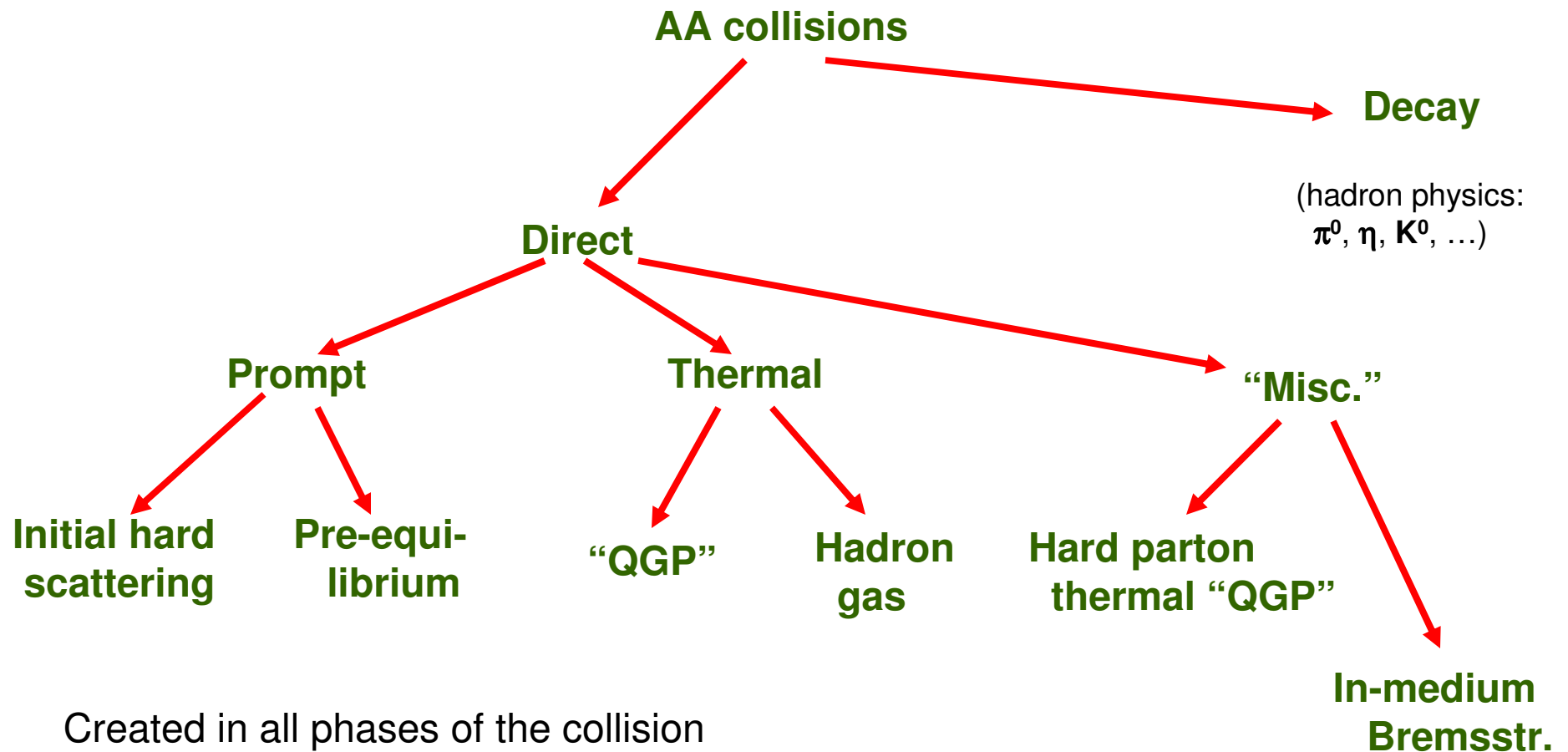
I will not go into details of how inclusive, direct, isolated, *etc.* photons are measured (experimental techniques discussed in detail in previous PHENIX Focus talks)

Relatively little about PHENIX



# Photons and photons

(A frequently used, albeit not unique terminology and hierarchy)



Created in all phases of the collision

Once created, they survive ( $\alpha_e \ll \alpha_s$ )  $\rightarrow$  time, temperature ... history

But this also makes measurements hard to interpret

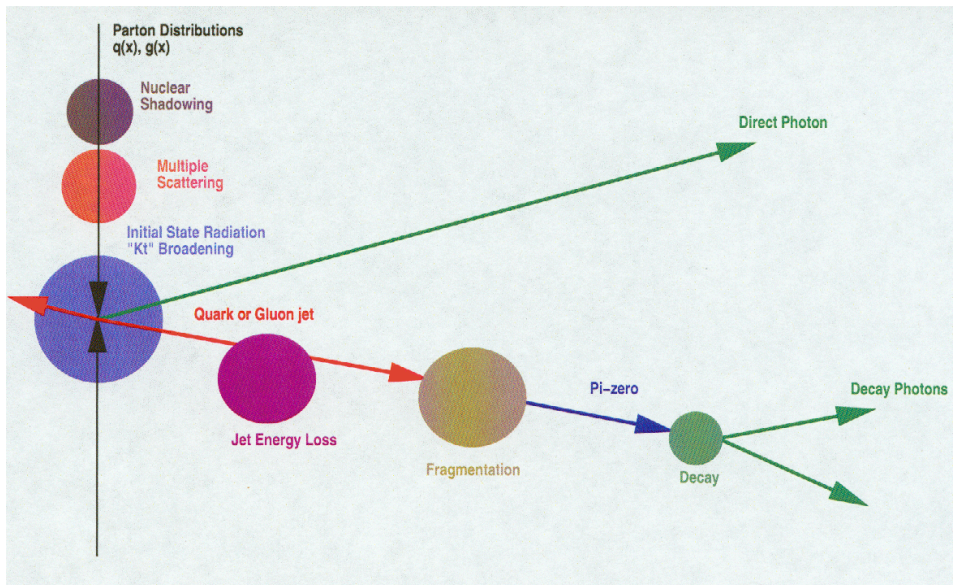
# Photons and collision systems



(Paul Stankus, ~1999)

A beautiful summary of the basic processes and how comparison of different signals in different colliding systems can **discriminate** between them

Some more details will be added later



	Parton Distributions $q(x), g(x)$	Nuclear Shadowing	Multiple Scattering	Initial State Radiation "Kt" Broadening	Jet Energy Loss	Fragmentation	Decay
<b>A+A</b>							
Inclusive Photons	✓	✓	✓	✓	✓	✓	✓
Inclusive Pi-Zero	✓	✓	✓	✓	✓	✓	
Direct Photons	✓	✓	✓	✓			
Photon-Pi-Zero Back to Back			✓	✓	✓	✓	
Photon-Photon Back to Back			✓	✓			
<b>p+p</b>							
Direct Photons	✓			✓			
<b>p+A</b>							
Direct Photons	✓	✓	✓	✓			

# Photons and their momenta



Many different components known / measured / to be disentangled

“Thermal window”:

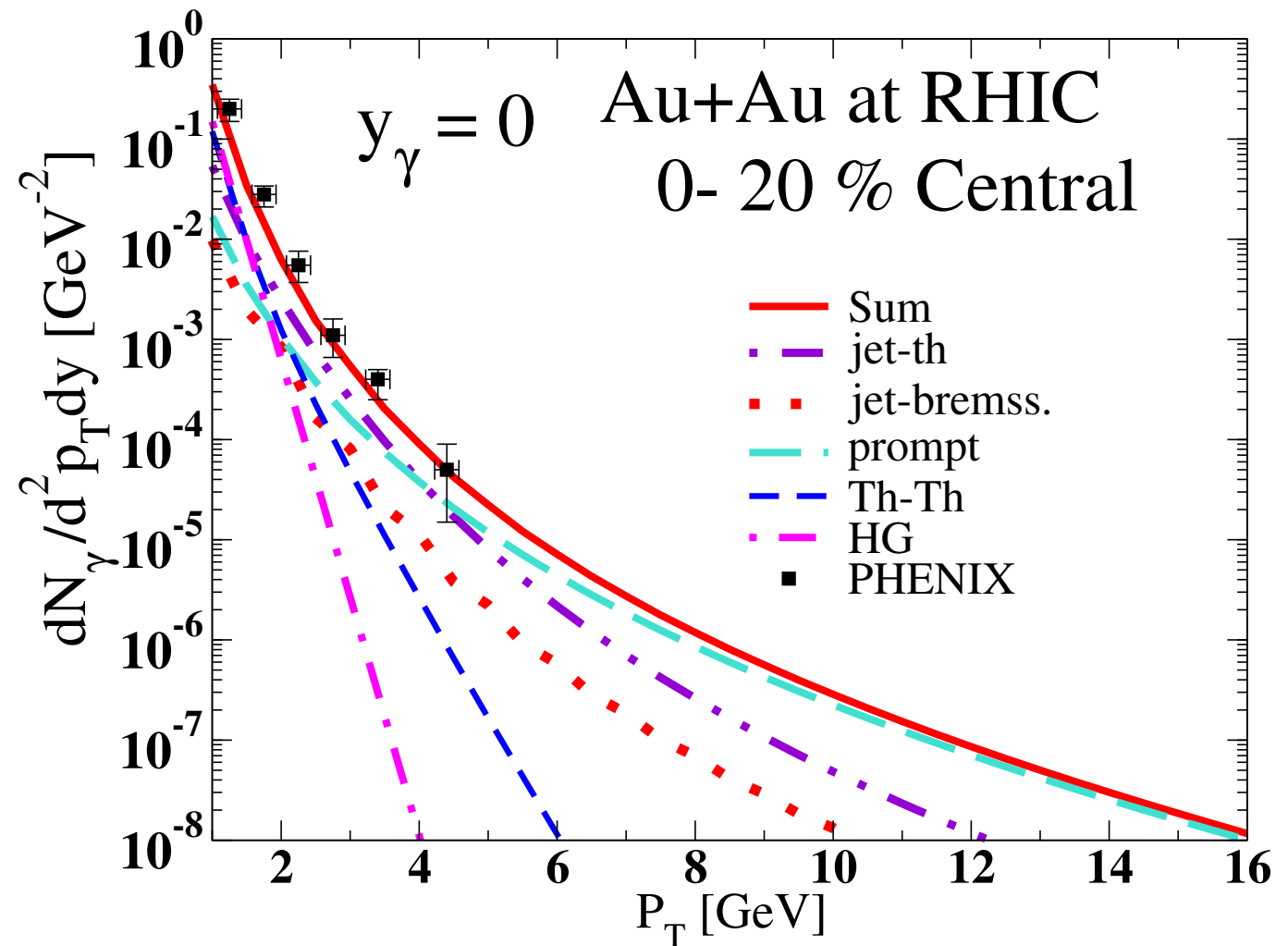
1-3 GeV (?)

“Jet-thermal”

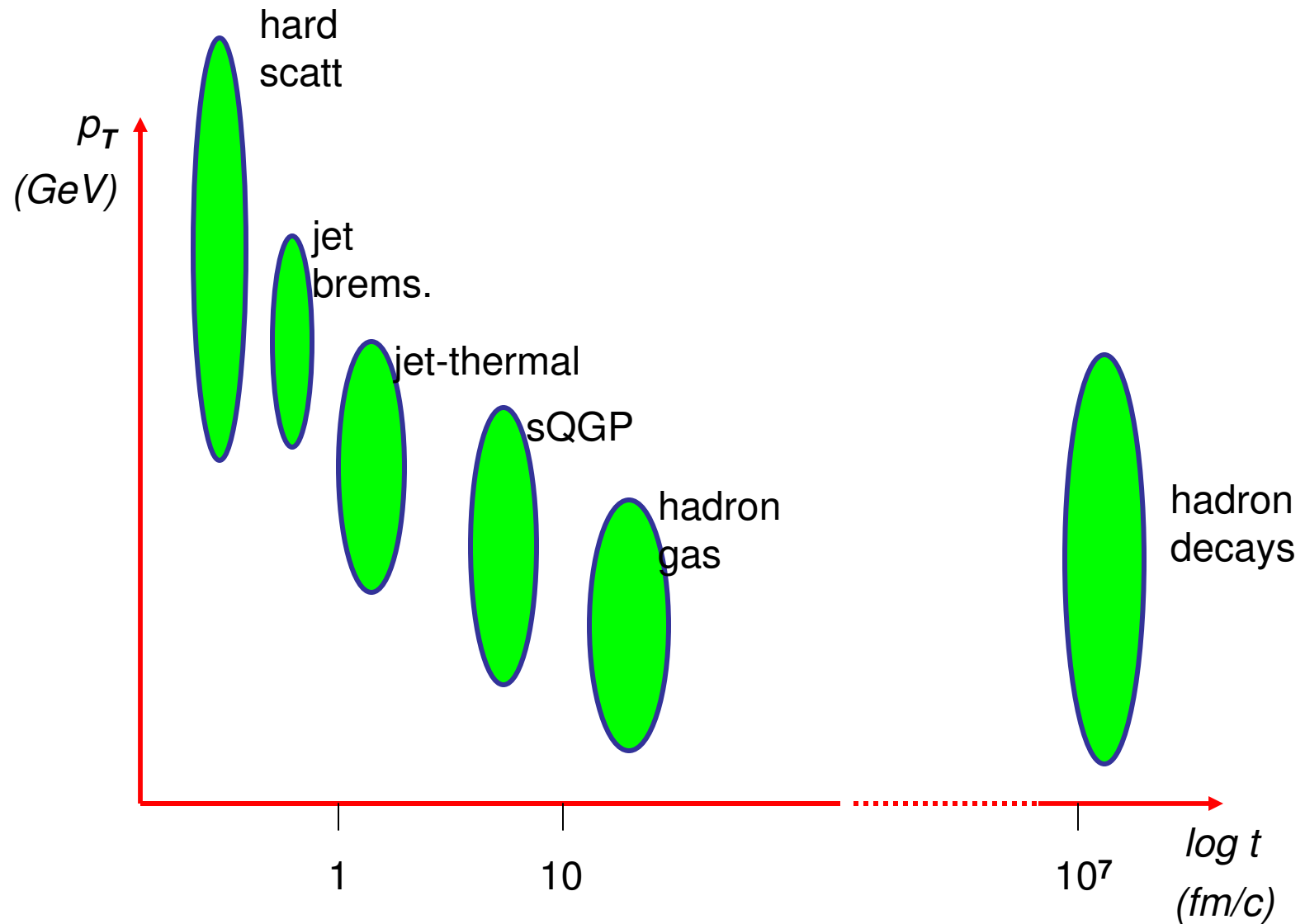
3-5 GeV (?)

Hard scattering

>6-7 GeV (?)



# Photons and time history of the collision





## Direct photons: the ways to present data

For identified particles usually their spectra are shown

- but for (**inclusive!**) photons this quantity is essentially **meaningless**
- direct photons are the (often small) difference of inclusive – hadron decay  $\gamma$   
→ often large errors, particularly at lower  $p_T$

Alternative ways of presenting the data (usually smaller syst. errors):

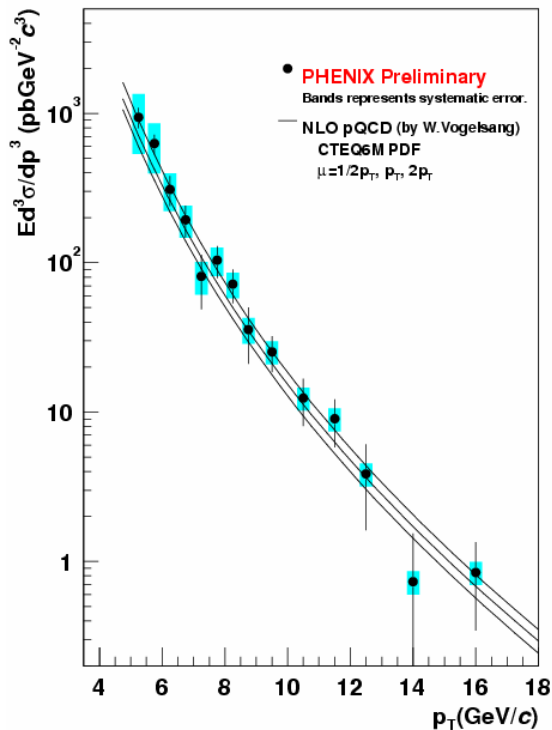
- $\gamma(p_T)/\pi^0(p_T)$  – often compared to the expected (simulated)  $\gamma/\pi$  under the assumption that all  $\gamma$  come from hadron decays  
→ deviation from the simulated  $\gamma/\pi$  indicates direct photons
- “double ratio”: the measured  $\gamma/\pi$  divided by the expected (hadronic only)  $\gamma/\pi$   
→ excess over 1 indicates direct photons

Hadron suppression in heavy ion collisions helps a lot!

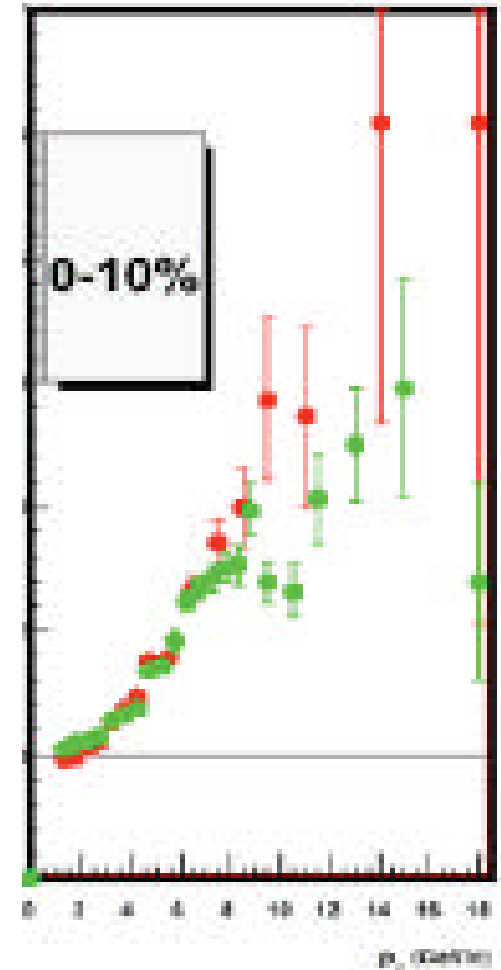
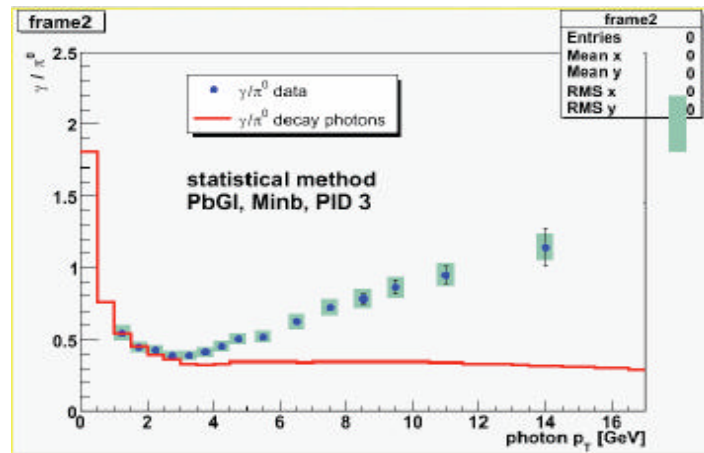


# Direct photons: spectrum, $\gamma/\pi$ and the “double ratio”

Complete information,  
but many syst. errors



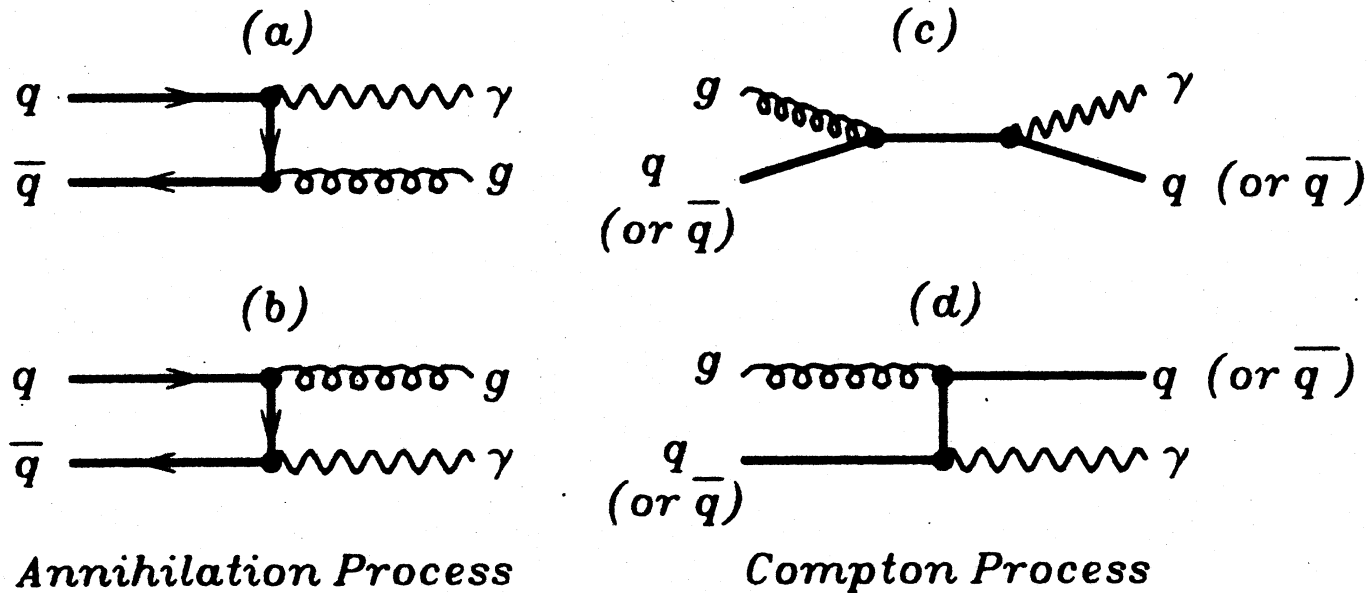
Less information, but more  
robust w.r.t. errors.  
These two presentations carry  
the same information







# Annihilation and Compton-scattering



Mandelstam  
variables:

$$s = (p_1 + p_2)^2 = (p_q + p_{\bar{q}})^2,$$

$$t = (p_1 - p_3)^2 = (p_q - p_\gamma)^2,$$

$$u = (p_1 - p_4)^2 = (p_2 - p_3)^2 = (p_{\bar{q}} - p_\gamma)^2.$$

# Annihilation – direction of the photon



Cross-section, averaged over quark spins, photon, gluon polarization (allow for mass):

$$\frac{d\bar{\sigma}}{dt}(q\bar{q} \rightarrow \gamma g) = \left(\frac{e_q}{e}\right)^2 \frac{8\pi\alpha_s\alpha_e}{s(s-4m^2)} \left\{ \left(\frac{m^2}{t-m^2} + \frac{m^2}{u-m^2}\right)^2 + \left(\frac{m^2}{t-m^2} + \frac{m^2}{u-m^2}\right) - \frac{1}{4} \left(\frac{t-m^2}{u-m^2} + \frac{u-m^2}{t-m^2}\right) \right\}.$$

which peaks when the denominators are minimum; writing them out explicitly gives

$$\begin{aligned} t - m^2 &= -2p_\gamma \cdot p_q \\ &= -2E_\gamma(E_q - |\mathbf{p}_q| \cos \theta_{\gamma q}). \end{aligned} \quad \rightarrow$$

$$u - m^2 = -2E_\gamma(E_{\bar{q}} - |\mathbf{p}_{\bar{q}}| \cos \theta_{\gamma \bar{q}}).$$

Both are minimum if  $\theta = 0$ ,  
i.e.  $\gamma$  lines up with either the  
quark or the antiquark:  
they are **colinear**

Expanding  $(t-m^2)^{-1}$  around  $\theta = 0$

$$\begin{aligned} \downarrow \quad \frac{1}{t-m^2} &\approx \frac{1}{-2E_\gamma[E_q - |\mathbf{p}_q| + |\mathbf{p}_q|\frac{1}{2}(\theta_{\gamma q})^2]} \\ &\approx -\frac{1}{E_\gamma E_q \left[ \left(\frac{m}{E_q}\right)^2 + (\theta_{\gamma q})^2 \right]}. \end{aligned}$$

giving the

approximate cone size

$$\rightarrow \Delta\theta_{\gamma q} = \frac{m}{E_q}$$



## Annihilation – energy of the photon

We've established that the most likely photon direction is collinear with the quark

The photon energy (one half of CMS energy) is

$$p_\gamma \cdot p_q + p_\gamma \cdot p_{\bar{q}} - s/2 = 0.$$

where  $p_\gamma \cdot p_q = E_\gamma(E_q - |\mathbf{p}_q| \cos \theta_{\gamma q}).$

In case of collinearity 
$$p_\gamma \cdot p_q \approx \frac{E_\gamma}{2E_q} m^2,$$

so after expanding  $\mathbf{s}$  and some approximations

$$p_\gamma \cdot (p_q + p_{\bar{q}}) - s/2 \approx (p_\gamma - p_q) \cdot p_{\bar{q}} + \frac{E_\gamma}{2E_q} m^2 - m^2.$$

Which in case of small  $\mathbf{m}$  (relativistic case) leads to 
$$\mathbf{p}_\gamma \approx \mathbf{p}_q,$$



# Compton-scattering

For Compton-scattering the argument is very similar to the annihilation, except that there is only one pole, in the **u** channel

The net result: the photon carries away most of the quark energy and flies out in a small cone around the original quark direction

Important note:

- the basic processes are the same, whether it is hard scattering or QGP; the difference is in the relevant parton distribution functions!
- for instance annihilation is a major contributor in QGP, but almost non-existent in p+p, since antiquarks are few and at small x
- so p+p “separates out” Compton-scattering → direct photons sensitive to the (polarized) gluon structure functions
- actually, this is not the entire story 😊

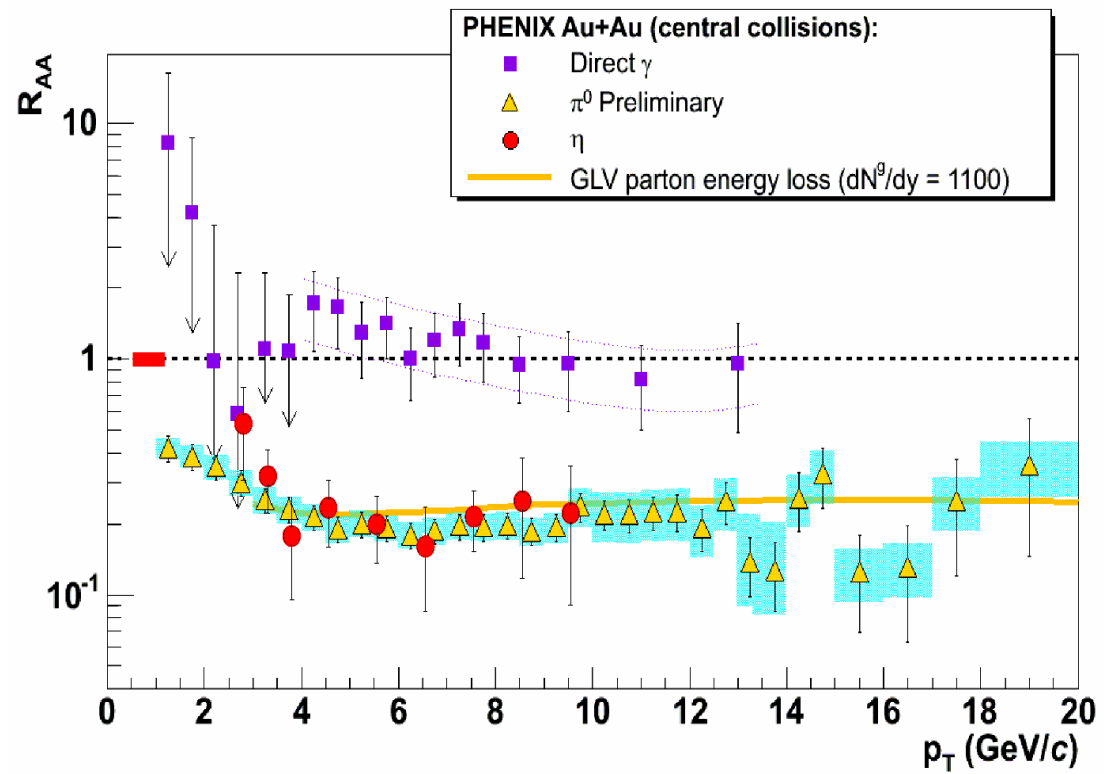
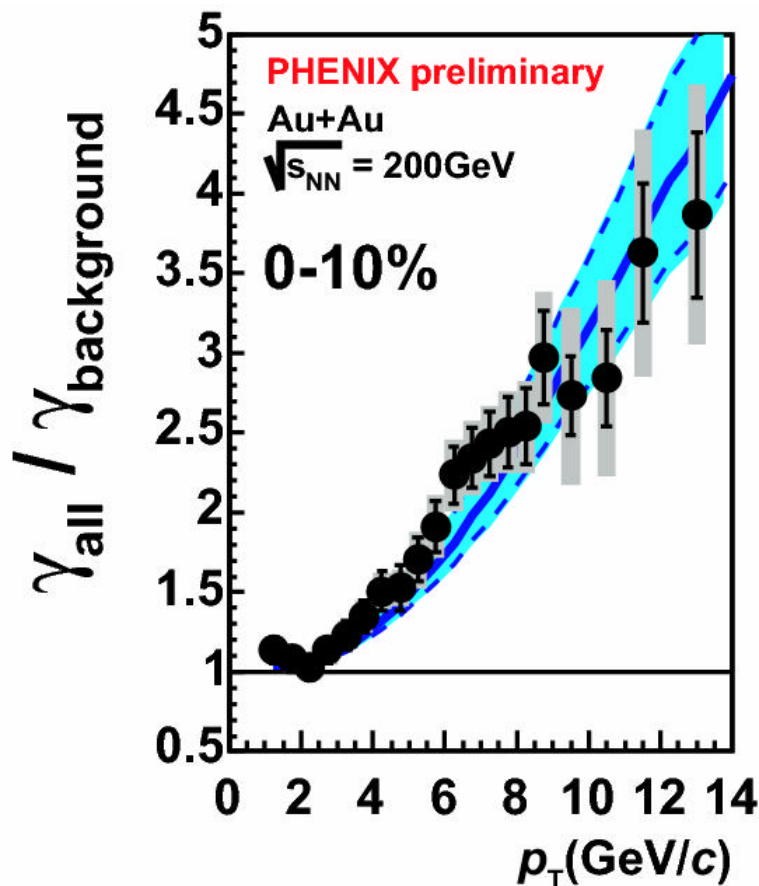
# Photons validating pQCD in heavy ion collisions



Remember: even if we did from the beginning  $N_{\text{coll}}$  scaling of pp to establish hadron suppression in AuAu, even if we did not see suppression in dAu, it wasn't proven beyond any doubt that suppression is due to the medium.

To prove it we needed something that does not interact with the medium, and actually does scale with  $N_{\text{coll}}$  - we needed photons

Note that this is also a statement on (non)modification of PDFs



# Where can pQCD be tested?

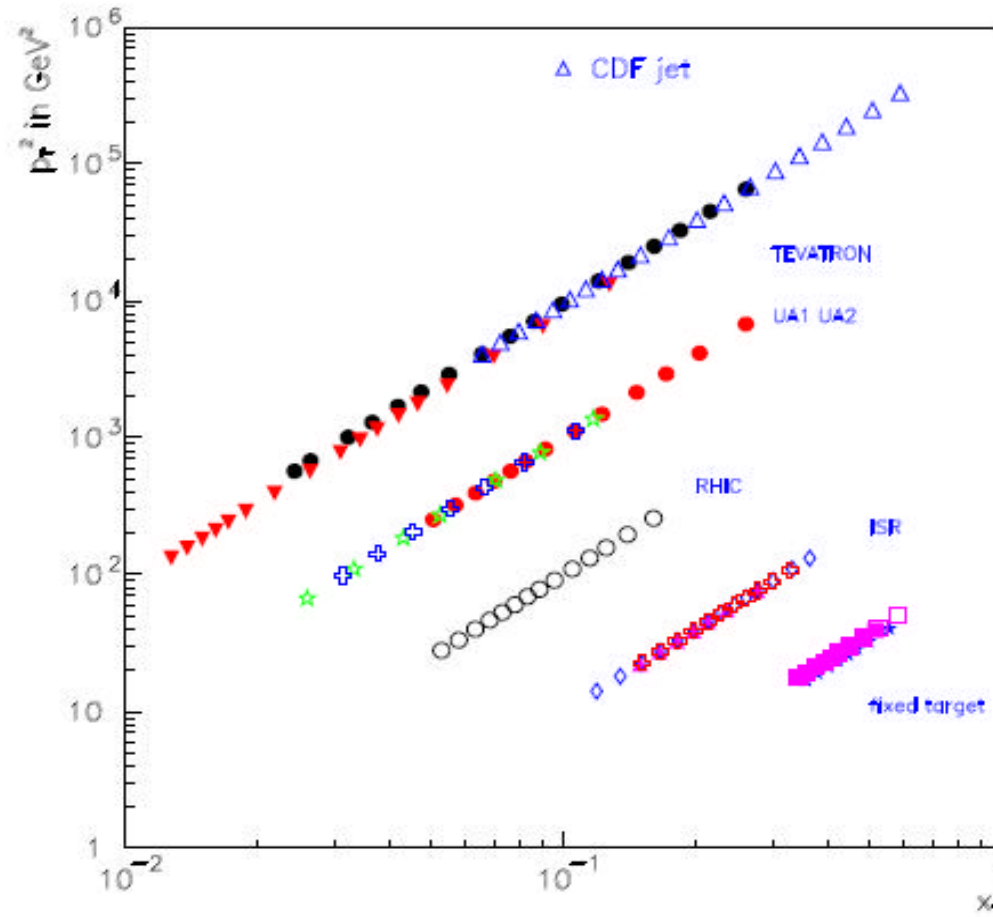


Figure 12: The kinematical region probed by prompt photon experiments compared to that relevant for jet production. Each data point is represented by a symbol as in Fig. 11 for photons, and by open triangles for jets.

hep-ph/0602133  
Aurenche et al.

# Test of pQCD: world data on pp vs theory



NLO calculations agree with data over 9 orders of magnitude – except for one experiment

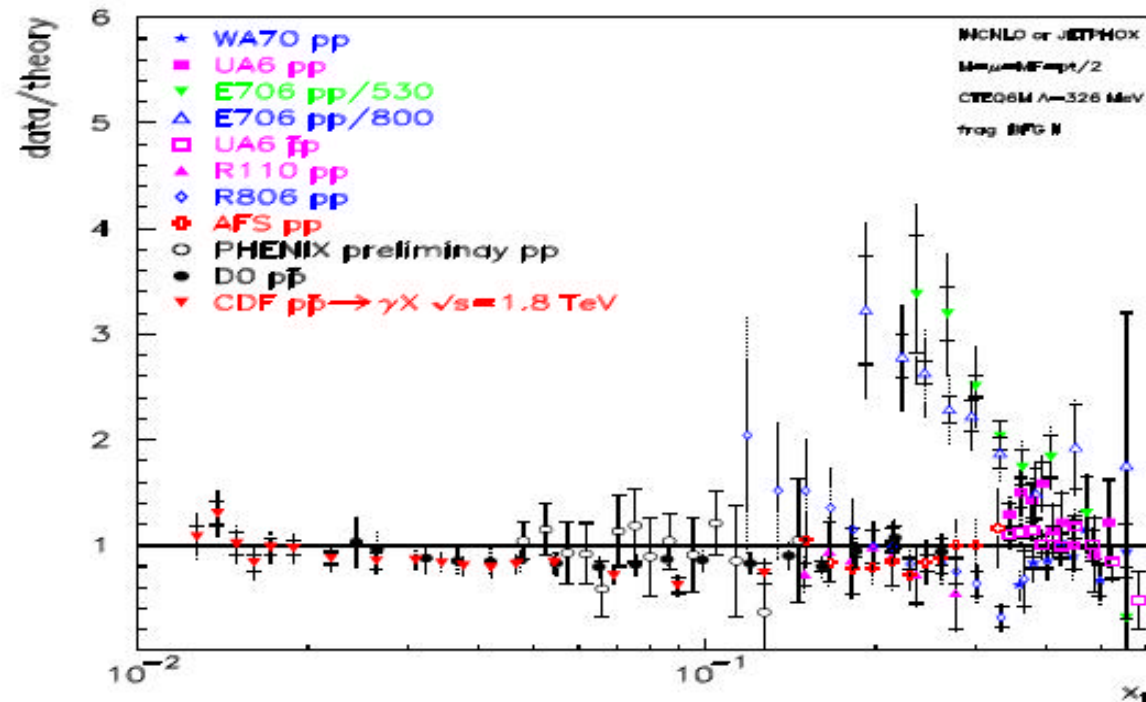


Figure 6: Ratios data/theory for collider and fixed target data with the scale  $\mu = p_T/2$ . For PHENIX and lower energy data the inclusive cross section is used while the isolated one is used for CDF and D0. Statistical errors only for PHENIX data.



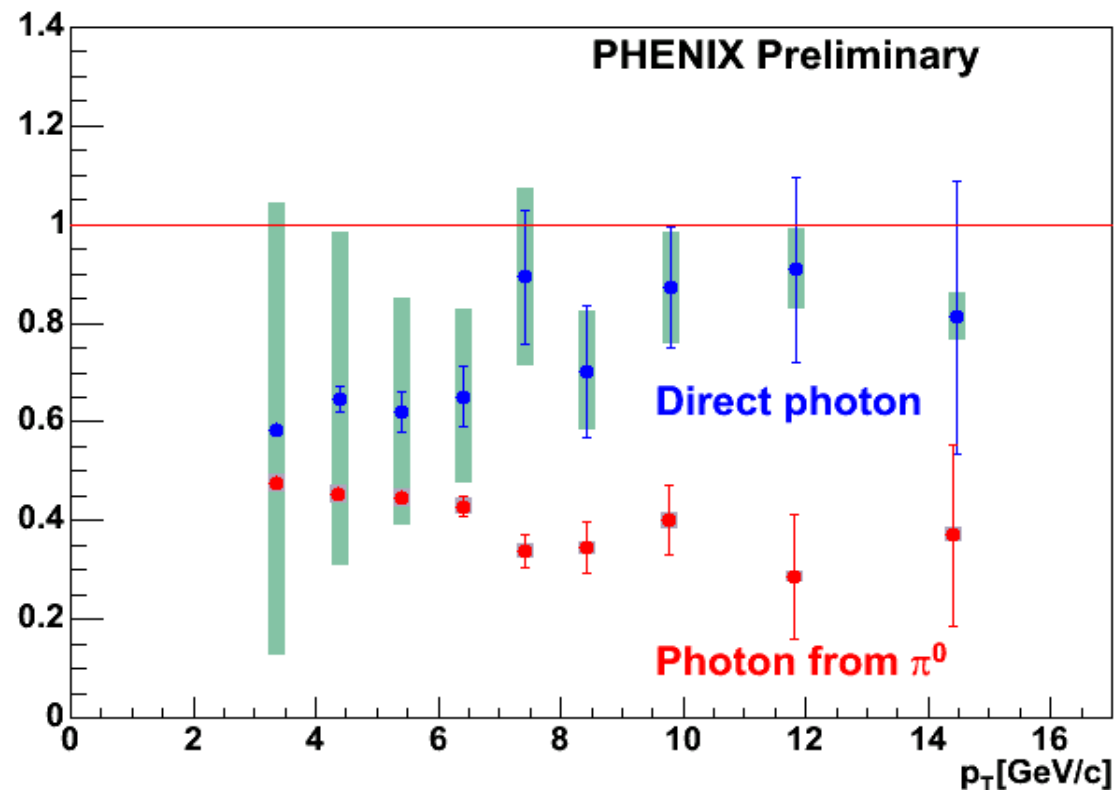
## Isolated / non-isolated photons

Fragmentation or  
hadronic photons  
should be accompanied  
by other particles nearby

Primordial (annihilation,  
Compton) photons should  
be isolated

Direct photon : isolation / subtraction

Photon from  $\pi^0$  : isolated photon / all



Hisa's CIPANP talk



# Isolated photons: experiment vs. theory

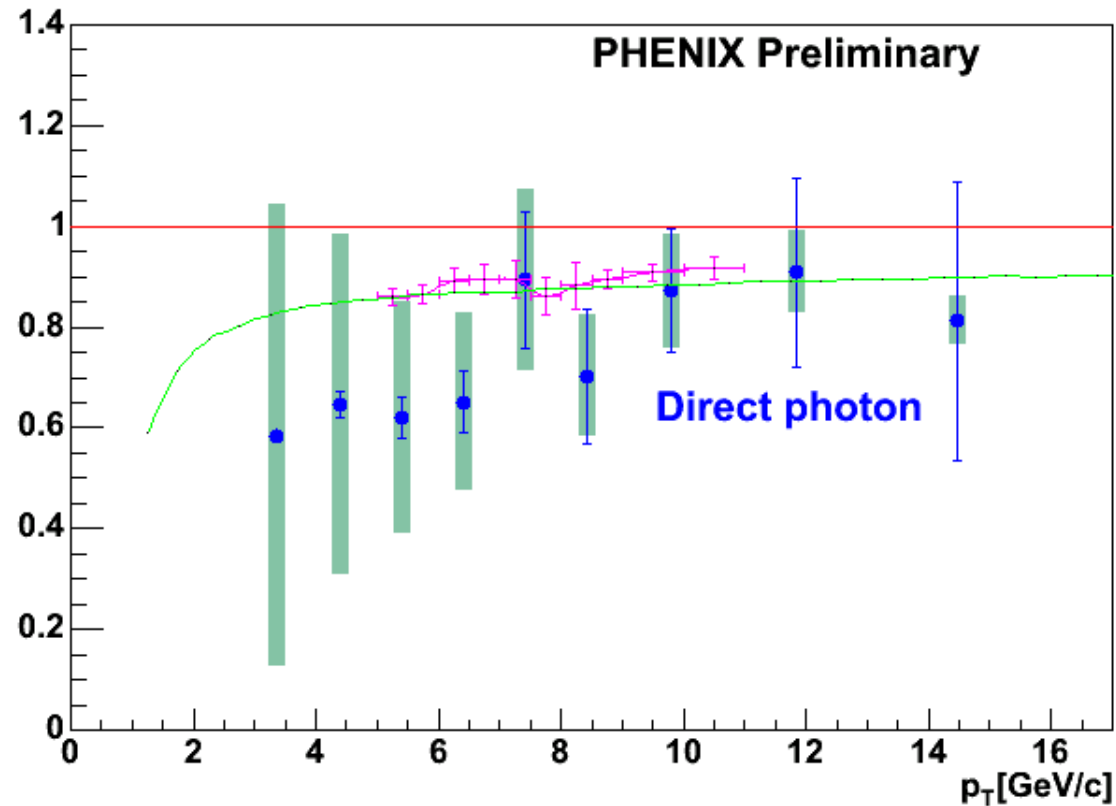


Isolation cut

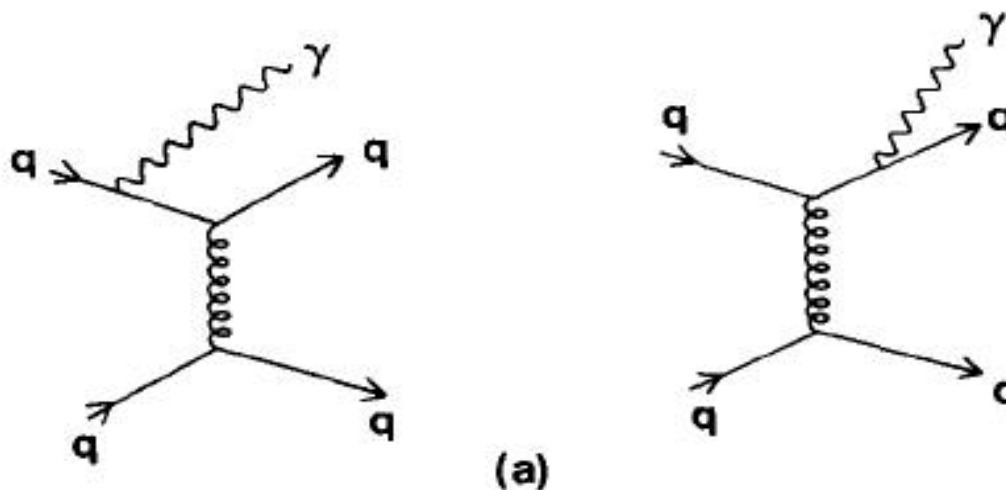
$$0.1 \cdot E_\gamma > E_{\text{cone}}(R=0.5\text{rad})$$

+ By M. Werlen,  
JETPHOX  
 $-.35 < y < .35$   
 $\mu = p_T$   
BFG set2, CTEQ6M

— By W. Vogelsang,  
 $R=0.4$   
 $\mu = p_T$ , CTEQ6M



# Bremsstrahlung – before and after



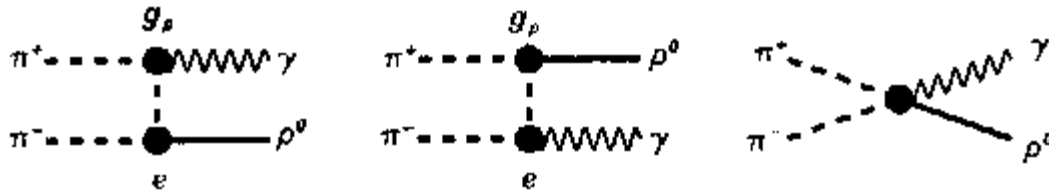
Before the interaction:

- modification of the effective PDF
- adds to the initial  $k_T$

After the interaction:

- photons in the vicinity of the leading hadron
- enhanced in the spatially large direction (out-of-plane)  $\rightarrow$  negative  $v_2$
- modifies effective FF

# Photon production from hadrons



$$\frac{d\sigma}{dt}(\pi^+\pi^-\rightarrow\gamma\rho^0) = \frac{8\pi\alpha_e\alpha_\rho}{s(s-4m_\pi^2)} \left\{ 2 - (m_\rho^2 - 4m_\pi^2) \left[ \frac{s-2m_\pi^2}{s-m_\rho^2} \left( \frac{1}{t-m_\pi^2} + \frac{1}{u-m_\pi^2} \right) + \frac{m_\pi^2}{(t-m_\pi^2)^2} + \frac{m_\pi^2}{(u-m_\pi^2)^2} \right] \right\}. \quad (16.69)$$

$$E_\gamma \frac{d\sigma}{dp_\gamma}(\pi^+\pi^-\rightarrow\gamma\rho^0) \approx \frac{1}{2}\sigma_{\pi^+\pi^-\rightarrow\gamma\rho^0}(s)E_\gamma[\delta(\mathbf{p}_\gamma - \mathbf{p}_{\pi^+}) + \delta(\mathbf{p}_\gamma - \mathbf{p}_{\pi^-})].$$

Bottom line (Lichard, ...): yield from hadron gas similar, maybe even bigger than from QGP (“the hadron gas outshines the plasma”)  
Effective photons temperature very high



# Jet-photon conversion

PRL 90 132301 (2003)  
(Fries, Muller)

Recall Compton-scattering:

the most probable process is that  
the photon is emitted in  $\sim$  the  
direction of the quark and carries  
 $\sim$  the entire energy of the quark  
(similar story for annihilation)

In the medium a jet (anti)quark can  
convert to a photon by annihilating  
or Compton scattering with a thermal  
gluon

Normal jets “encode” energy loss in the  
entire plasma, jet-conversion photons  
have only part of that (harder spectrum  
than “ordinary” jets)

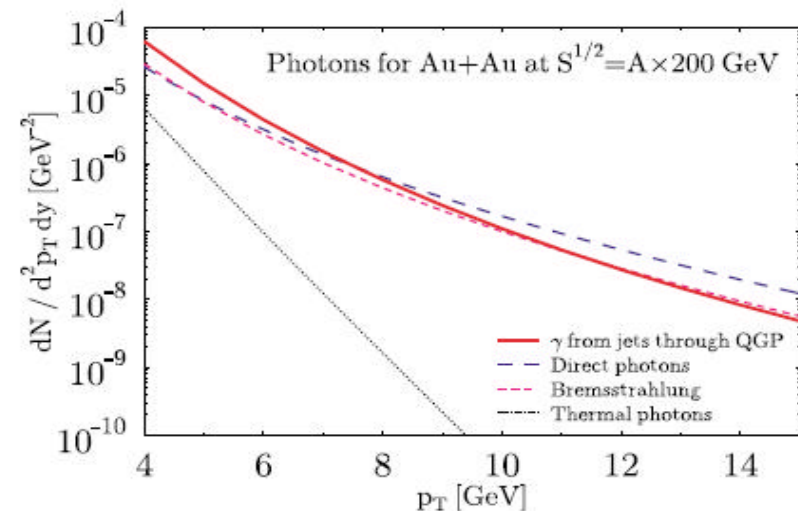


FIG. 1 (color online). Spectrum  $dN/d^2p_\perp dy$  of photons at  $y = 0$  for central collision of gold nuclei at  $\sqrt{s_{NN}} = 200$  GeV at RHIC. We show the photons from jets interacting with the medium (solid line), direct hard photons (long dashed), bremsstrahlung photons (short dashed), and thermal photons (dotted).

## Oh, I almost forgot: thermal radiation ☺

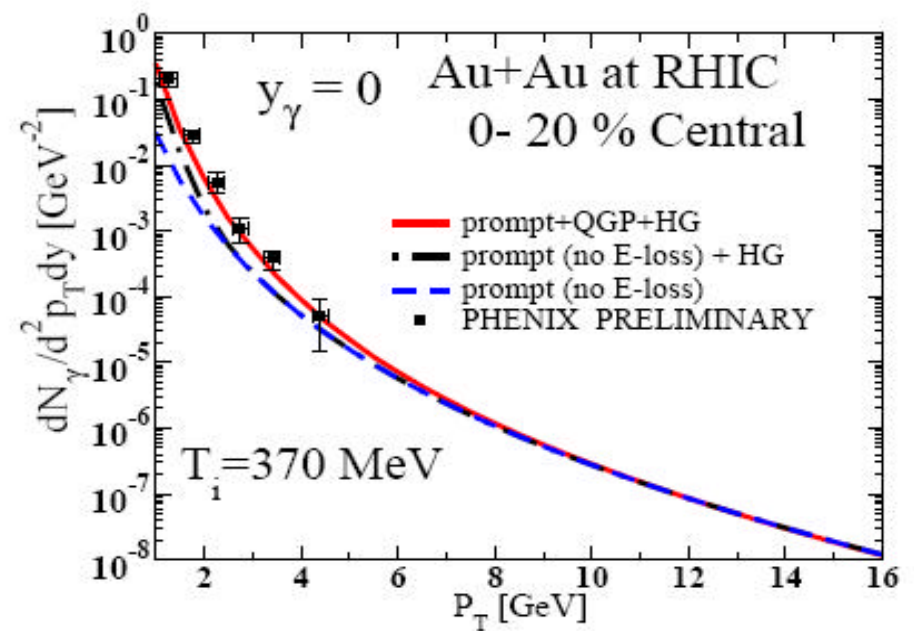
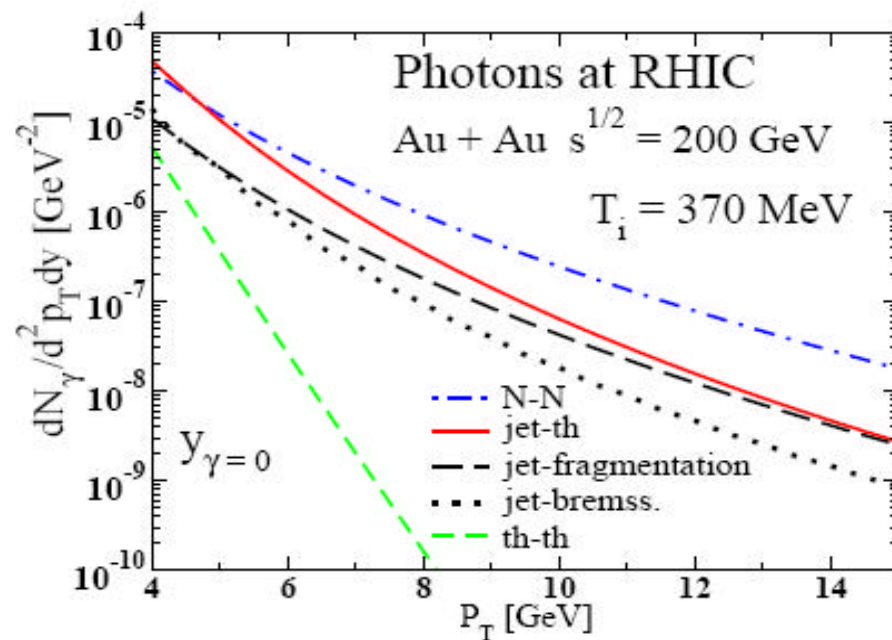


hep-ph/0512109 (Gale)

Once upon a time this **was** the issue most talked about...

Hard to access experimentally

Phase transition



# Annihilation vs. Compton (Fermilab – E706)



PRD 70 (2004) 092009

(Apanasevich et al.)

In  $\pi^-$  Be annihilation is significant due to the (-2/3 charge) valence anti-u while it would be suppressed in  $\pi^+$  Be since anti-d is 1/3 only

In pBe you only have sea antiquarks  
-> Compton dominates

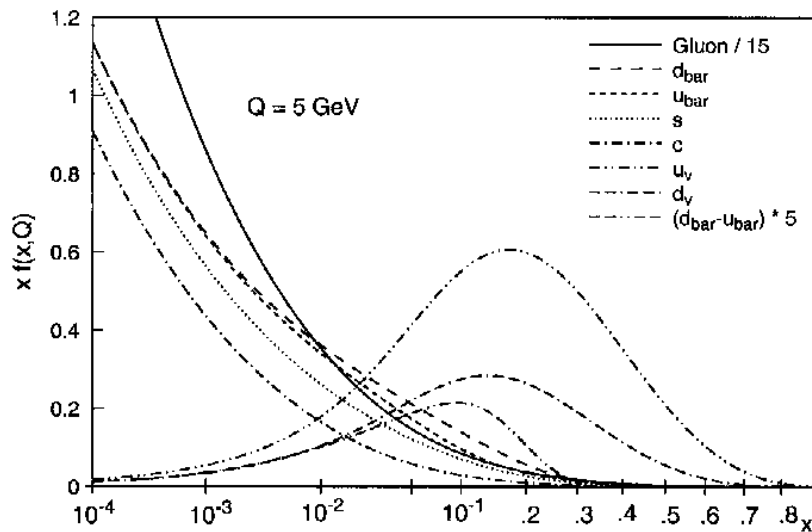


Figure 2 Overview of CTEQ5M parton distributions at  $Q = 5$  GeV. The gluon distribution is scaled down by a factor of 15, and the  $(\bar{d} - u)$  distribution is scaled

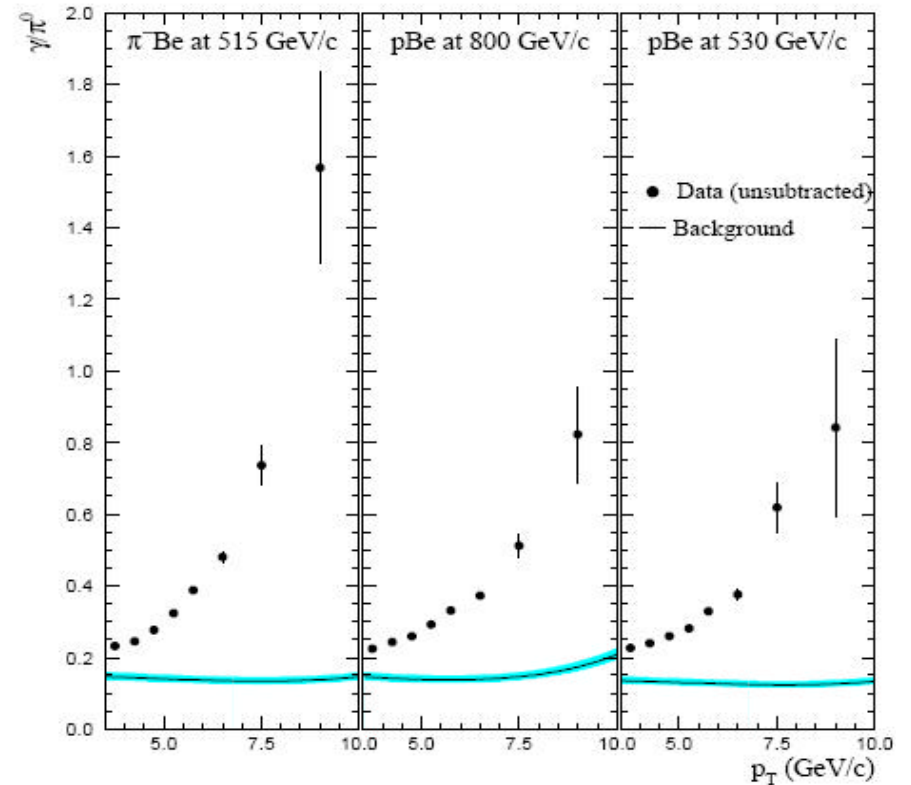


FIG. 21: Ratios of the 90N direct-photon candidate spectra to the measured  $\pi^0$  cross sections (points) compared to  $\gamma_{bkg}/\pi^0$  from the DGS (curves) as functions of  $p_T$  for the data samples considered in this paper. The error bars represent statistical contributions to the uncertainties. The width of each band around the background represents the systematic uncertainty on that background.

# Initial soft-gluon radiation? (Fermilab – E706)



PRD 70 (2004) 092009  
(Apanasevich et al.)

NLO pQCD (dashed curves) does not describe photon production  
Assuming that initial (pre-collision) soft gluon radiation introduces an effective  $k_T$ , the data are well described (solid curves)  
No other experiment needs such  $k_T$  enhancement (non-perturbative)

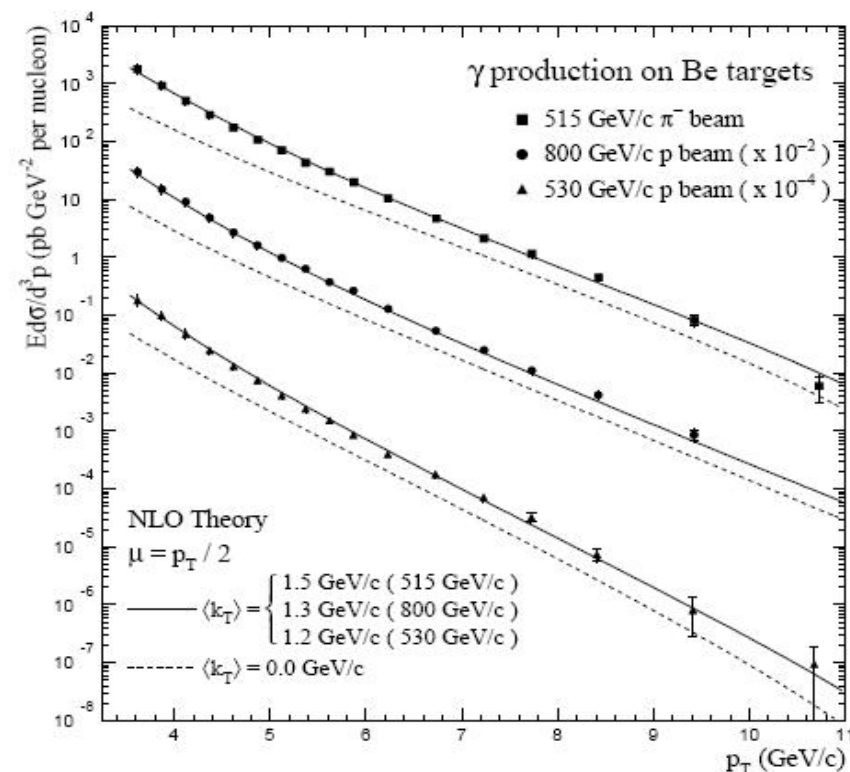


FIG. 28: Invariant differential cross sections per nucleon for direct-photon production as functions of  $p_T$ , averaged over rapidity, for 515 GeV/c  $\pi^-$  and 800 and 530 GeV/c proton beams incident upon beryllium. The error bars represent the statistical and systematic uncertainties combined in quadrature; the innermost interval indicates the statistical uncertainties. Overlaid on the data are NLO PQCD and  $k_T$ -enhanced NLO PQCD calculations. GRV92 PDF were used in the incident  $\pi^-$  calculations, while CTEQ5M PDF were used in the incident proton calculations.



# Cronin-effect in photons? (Fermilab – E706)



hep-ex/0506003  
(Apanasevich et al.)

Difference between  $\gamma$  (initial state effect only?) and  $\pi^0$ ;  
still if true, confusing picture

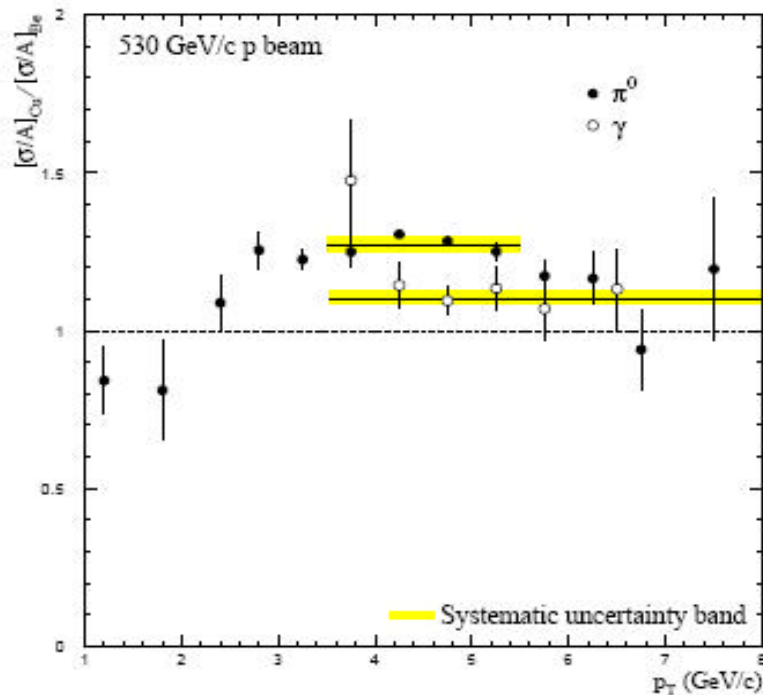


FIG. 16: The ratio of inclusive  $\pi^0$  and direct-photon production cross sections per nucleon in  $p\text{Cu}$  to those in  $p\text{Be}$  collisions at 530 GeV/c. Simple straight line fits to regions with relatively flat distributions have been overlaid on the data. The error bars represent only statistical contributions to the uncertainties. Systematic uncertainties are indicated by the shaded region associated with the fit.

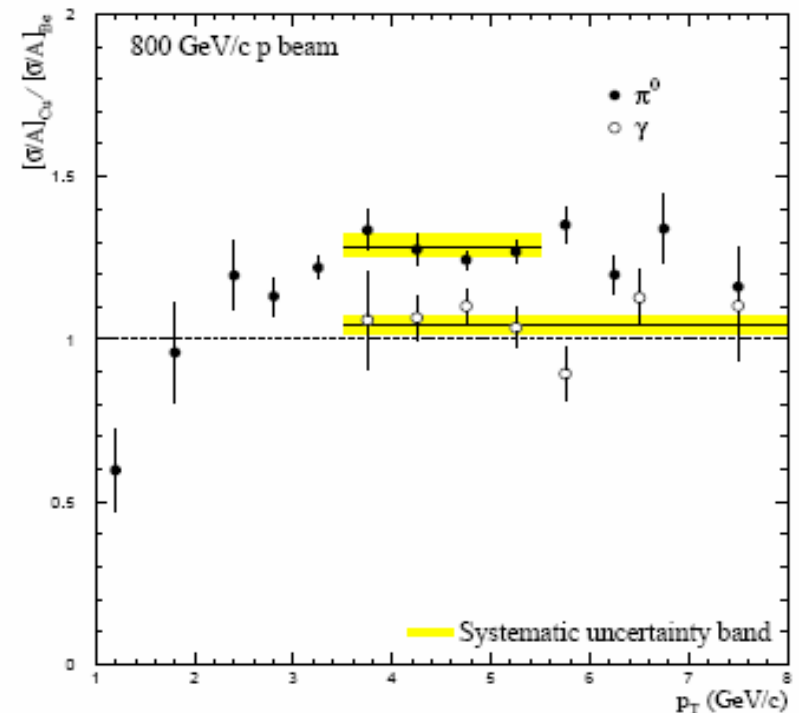


FIG. 17: The ratio of inclusive  $\pi^0$  and direct-photon production cross sections per nucleon in  $p\text{Cu}$  to those in  $p\text{Be}$  collisions at 800 GeV/c. Simple straight line fits to regions with relatively flat distributions have been overlaid on the data. The error bars represent only statistical contributions to the uncertainties. Systematic uncertainties are indicated by the shaded region associated with the fit.





## CERN – R110

Nucl. Phys. B327 (1989) 541  
(Angelis et. al, inc. MJT...)

pp at 63 GeV

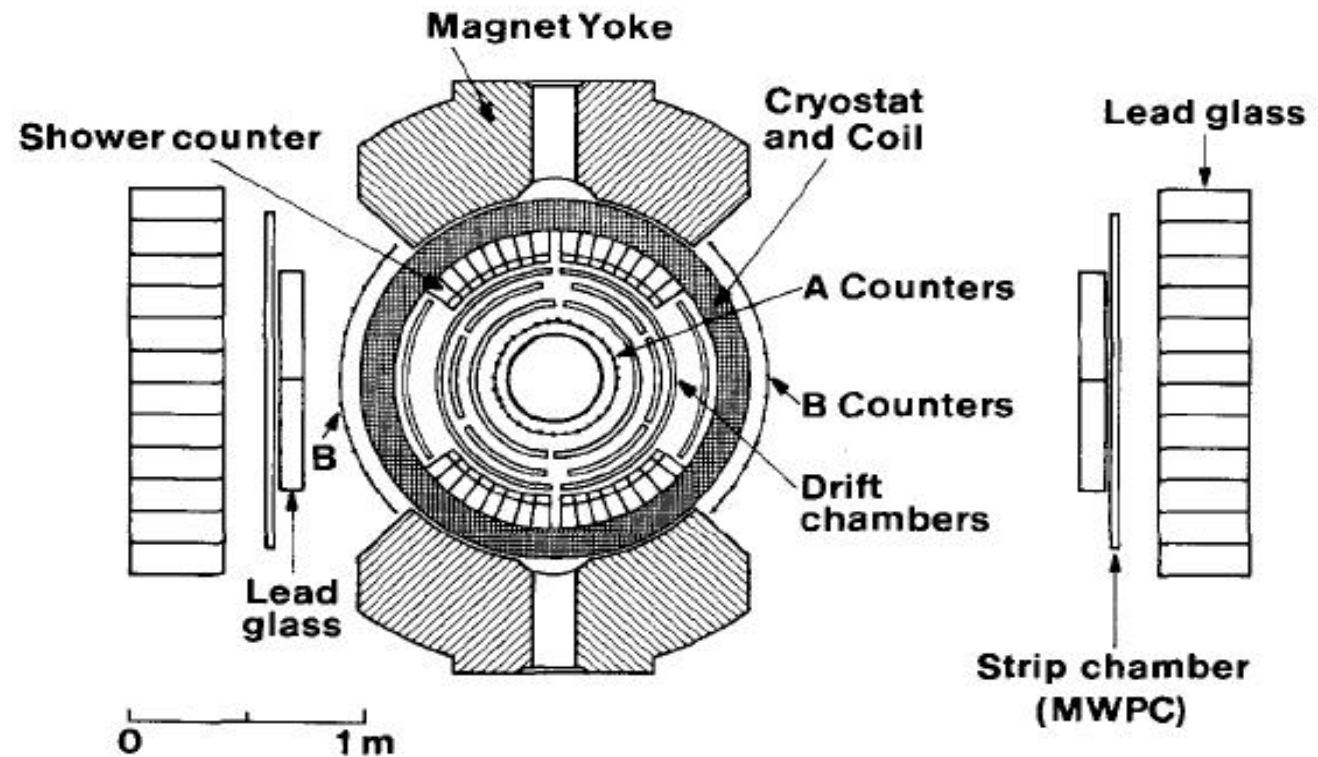


Fig. 2. The R-110 apparatus viewed along the beam axis.

(I couldn't resist showing this...)

# Isolated photons setting the jet E-scale



Nucl. Phys. B327 (1989) 541  
(Angelis et. al, inc. MJT...)

If in a back-to-back  $\gamma$ -jet pair  
the photon is isolated, it is  
a good measure of the total  
jet energy (modulo initial  $k_T$ )

The complete disappearance of  
 $\gamma$ -side partners with  $p_T > 1.0$  GeV  
also suggests that Bremsstrahlung  
is not a dominant source of photons  
here

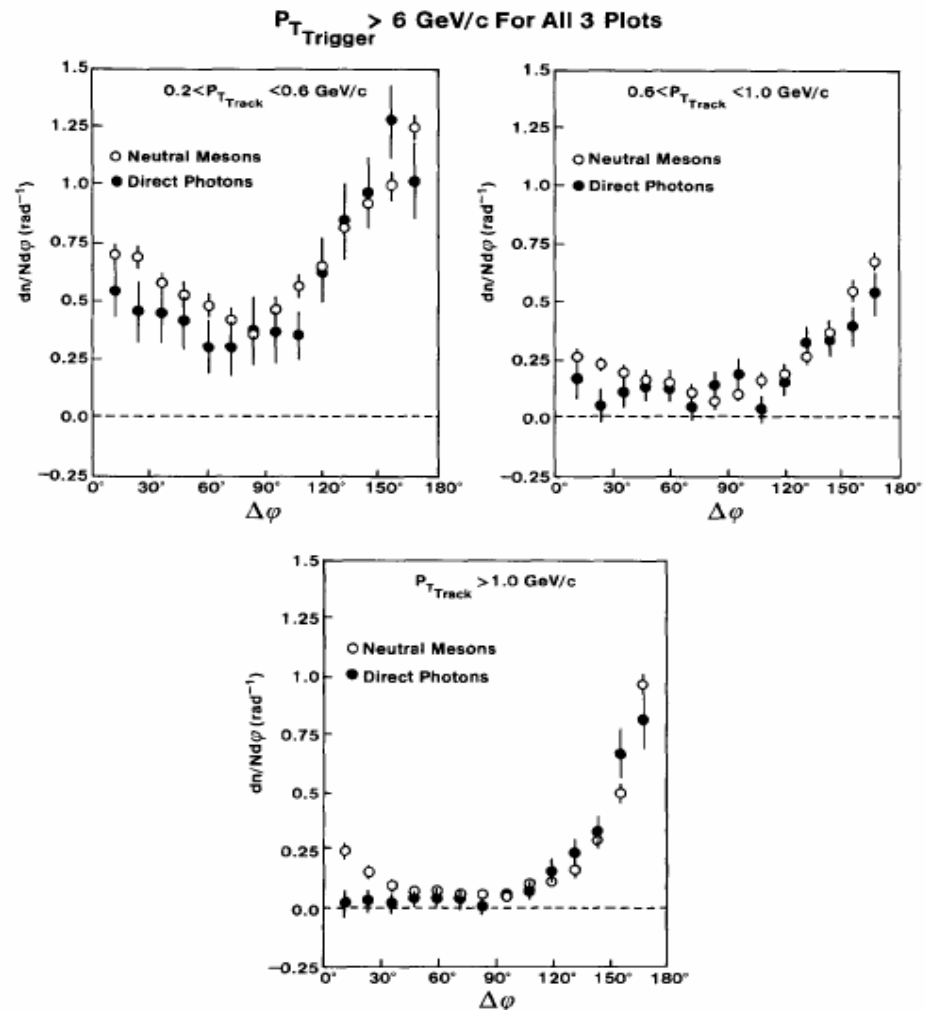


Fig. 15. Azimuthal distributions of charged associated particles for the direct photon and neutral meson samples (after background extraction) for  $p_{T,Trigger}$  greater than 6.0 GeV/c and for three different  $p_{T,Track}$  ranges.

## Which graphs dominate direct $\gamma$ production?



Another sign of the dominance of the Compton-graph:

- for Compton the recoil particle is a quark
- the proton is **uud** and **u** couples 4 times stronger to photons
- for photons the probability to get a positive hadron away-side is higher (except when we start submerging in the sea)

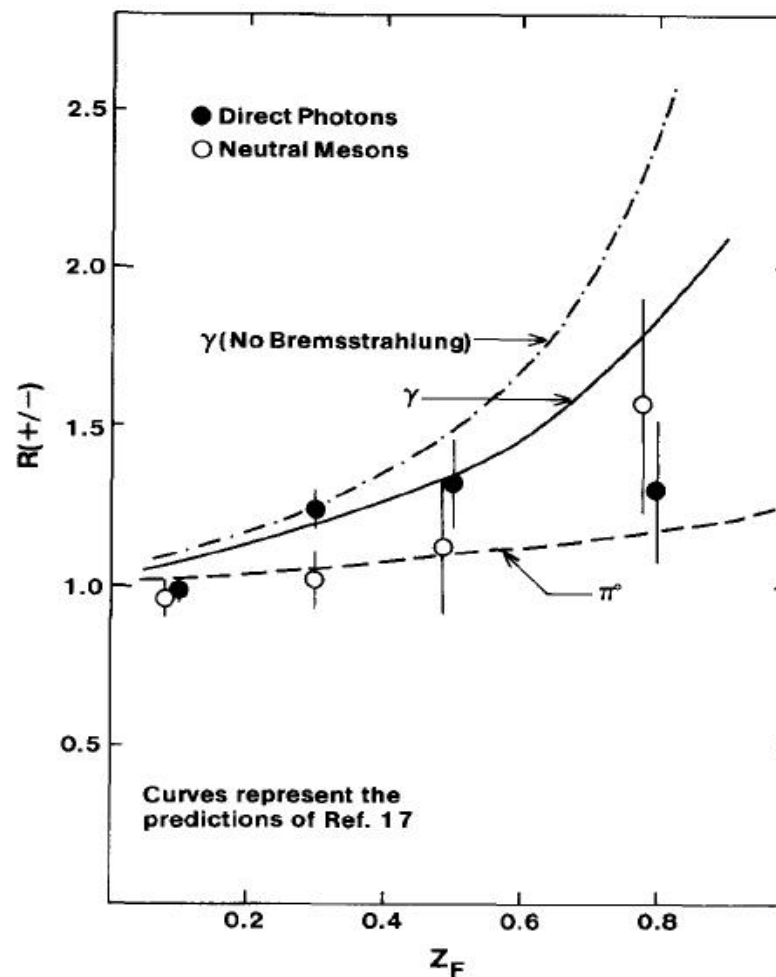


Fig. 16. Ratio of the number of positive to negative charged hadrons in the away-side vs.  $z_F$  for the direct photon and neutral meson samples. Also shown are the predictions of ref. [17].

# Photon flow - direct, hadronic, ...



WA98

Eur. Phys. J. C41 (2005) 287

Decreasing flow with increasing centrality

Photon  $v_2$  not plotted directly

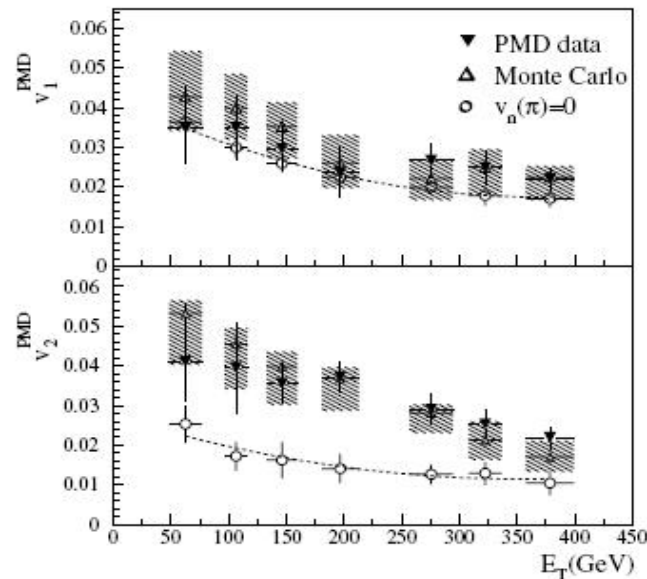


Fig. 6. a First order,  $v_1^{\text{PMD}}$ , and b second order,  $v_2^{\text{PMD}}$ , photon anisotropy coefficients in the pseudorapidity region  $3.25 \leq \eta \leq 3.75$  for different centralities are shown by filled triangles. Statistical and systematical errors are added in quadrature and shown as bars on the filled triangles. Open triangles are the most probable values of  $v_n^{\text{PMD}}$  as expected from the simulation. The shaded regions indicate the simulation uncertainties as described in Sect. 5.4. The open circles show the calculated values of  $v_n^{\text{PMD}}(v_n(\pi) = 0)$  assuming an isotropic distribution of pions with the dashed curve indicating a smooth polynomial fit to the open points. Note, however, that  $v_n^{\text{PMD}}(v_n(\pi) = 0)$  can not be directly subtracted from the  $v_n^{\text{PMD}}(v_n(\pi) > 0)$  to obtain the anisotropy flow coefficients  $v_n$ , as explained in the text

PHENIX PPG046

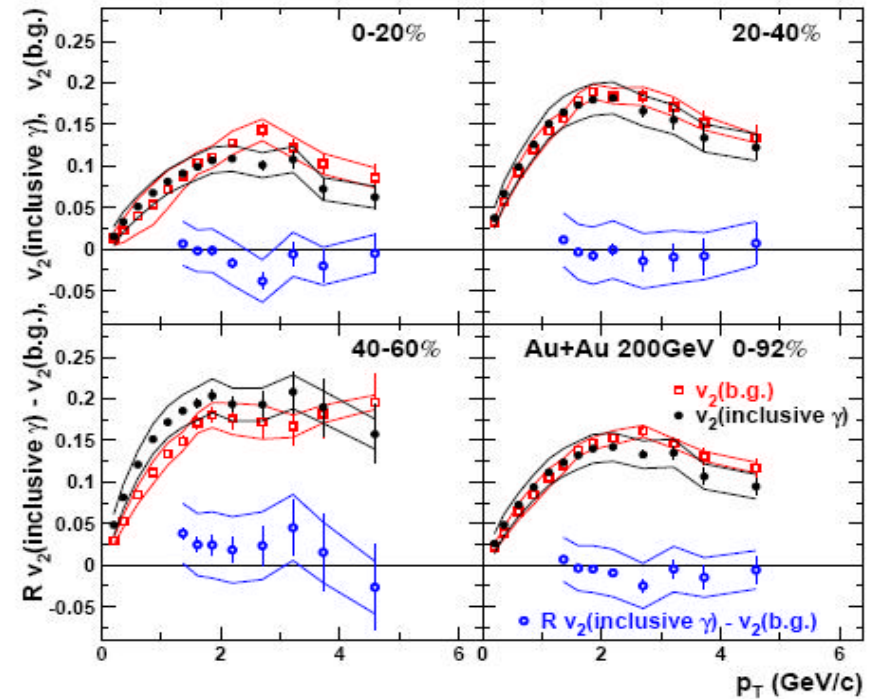


FIG. 2: (Color online) The measured  $v_2$  of inclusive photons ( $v_2^{\text{inclusive } \gamma}$ , solid circle) and expected photon  $v_2$  from hadronic decay ( $v_2^{\text{b.g.}}$ , open square). A subtracted  $v_2$  quantity  $Rv_2^{\text{inclusive } \gamma} - v_2^{\text{b.g.}}$  is plotted at the bottom of each panel (open circle), where  $R = (N_{\text{direct } \gamma} + N_{\text{b.g.}})/N_{\text{b.g.}}$ . The quantity corresponds to a product of the direct photon  $v_2$  and a positive factor  $R - 1$ , ( $v_2^{\text{direct } \gamma}(R - 1)$ ).

# Photon flow: the point



hep-ph/0508201  
(Turbide, Gale, Fries)

Primary hard scattering photons  
should be produced isotropically

Hadron decay photons should  
inherit the flow of hadrons  
(positive  $v_2$ )

If jet-plasma interaction is  
significant source of photons,  
negative (out-of-plane,  $\pi/2$  shift)  
flow is expected at high  
momenta (“inverse optical  
mechanism”, the plasma is  
thicker out-of-plane)

Note the  $p_T$  scale

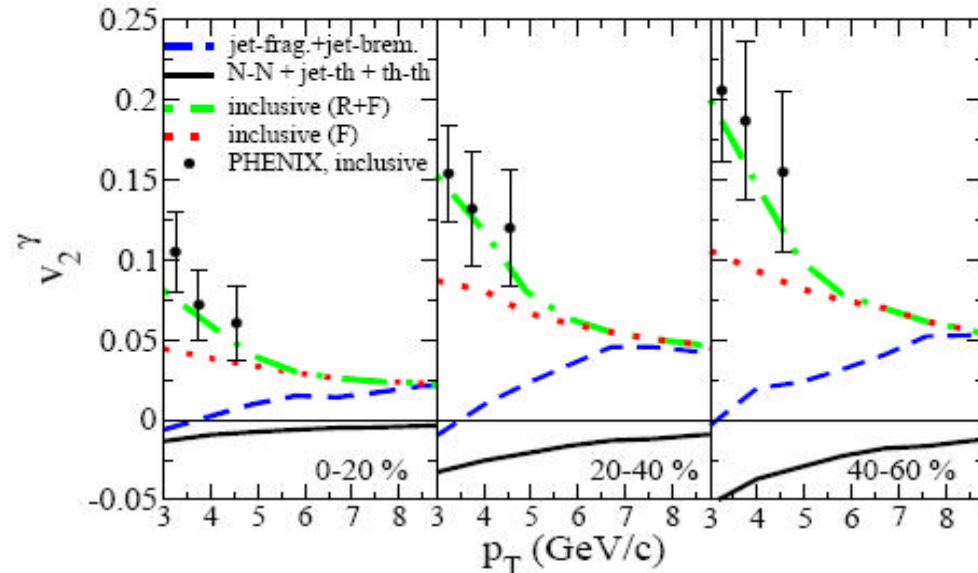


FIG. 2: (Color online)  $v_2$  as a function of  $p_T$  for Au+Au collisions at RHIC. The dashed line shows jet-fragmentation and induced bremsstrahlung only while the solid lines give jet-photon conversion, primary hard and thermal photons. The dotted line shows direct photons and the background from decay of neutral mesons coming from jets. The dot-dashed line adds photons from decay of recombined pions as well and can be compared to the inclusive photon  $v_2$  measured by PHENIX [20].



# Space-time evolution: photon HBT



Theory: photons only

Since  $\pi^0$  decays at  $\sim 10^7$  fm,  $\Delta q \sim$  few eV – no influence on measurable correlation

- although HBT corr. of the  $\pi^0$ -s themselves may introduce fake correlation

Correlations at various  $p_T$  can map out space-time evolution (in principle...)

$\Delta q \sim 1/R$   $p_T$  range  $\sim 1/t$

Exp:  $\pi^0$  turned on (WA98)

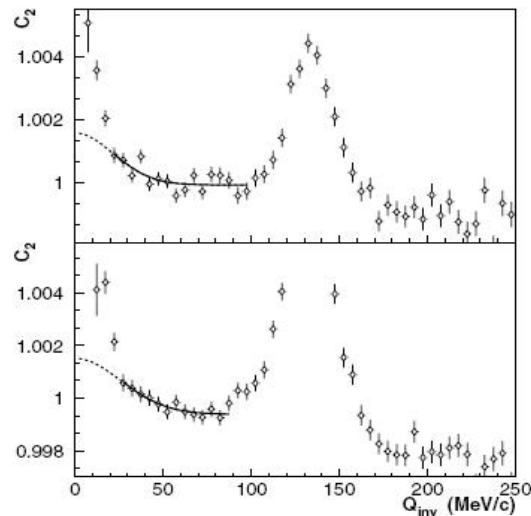


FIG. 1. The two-photon correlation function for narrow showers with  $L_{\min} > 20$  cm (diamonds) and average photon momenta  $100 < K_T < 200$  MeV/c (top) and  $200 < K_T < 300$  MeV/c (bottom) fitted with Eq. (1). The solid line shows the fit result in the fit region used (excluding the  $\pi^0$  peak at  $Q_{\text{inv}} \approx m_{\pi^0}$ ) and the dotted line shows the extrapolation into the low  $Q_{\text{inv}}$  region where backgrounds are large.

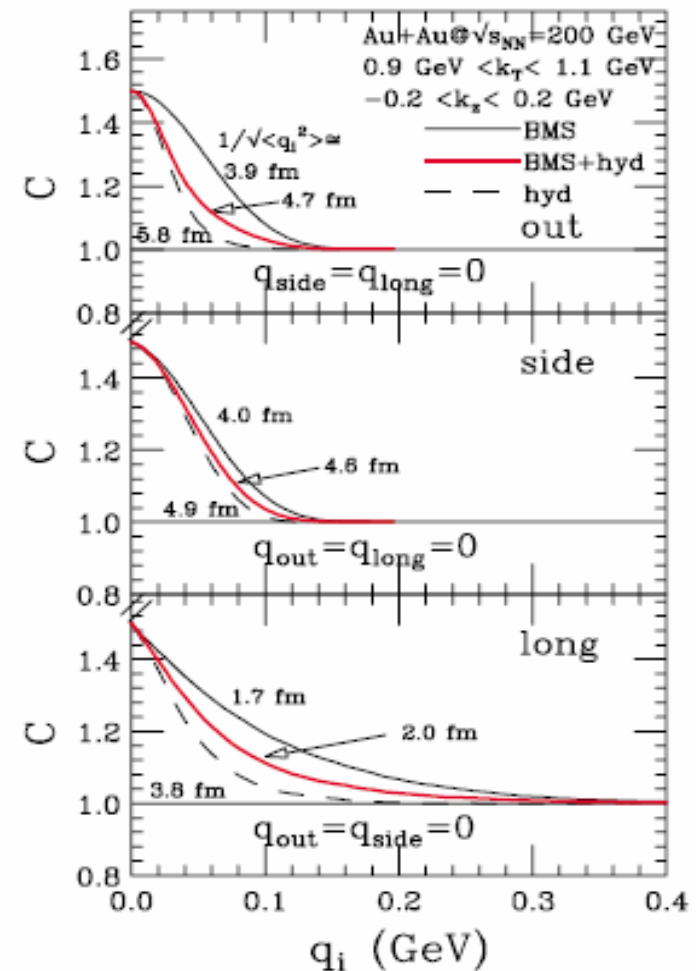


FIG. 3 (color online). Outward, sideward, and longitudinal intensity correlation of photons at 1 GeV, considering only PCM (BMS), only thermal (hyd), and all PCM + thermal photons (BMS + hyd).

## Summary: disentangling processes (?)



We started out saying that photons (penetrating probes!) report on every stage of the collision; but it's a mixed blessing – can we **disentangle** those processes?

Sometimes we can. Examples (very far from complete list):

- testing LO processes using suitable **hh** collisions
- set the jet energy scale in  $\gamma$ -jet pairs with the photons isolated
- access system size at earliest times (high  $p_T$  HBT)
- medium properties via photons in the vicinity of a leading hadron
- medium properties via (out-of-plane?) photon flow
- medium properties via E-loss in jets vs jet-photon conversion  $\gamma$
- temperature history of the collision (tall order!)
- access to gluon structure functions (when Compton dominates)
- ...

If you ask for a signal, that encodes everything – watch out!  
You may get what you prayed for...

